

Blocks in the Asymmetric Simple Exclusion Process

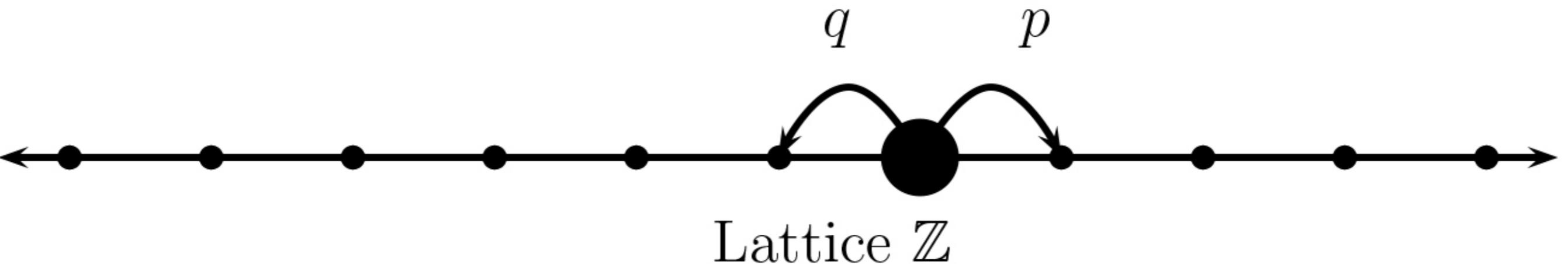
Painlevé Equations and Applications:
A Workshop in the Memory of A.A. Kapaev
Ann Arbor, August 2017

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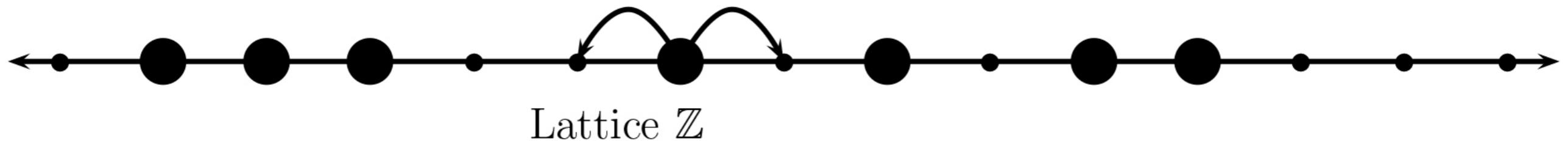
Simple Random Walk



- Can make **time continuous** by giving particle a “random alarm clock”, i.e. exponential distr. with mean 1.
- This is arguably one of the most important, if elementary, *stochastic processes*.
- Want **many particles**—to be interesting these particles must interact.

Asymmetric Simple Exclusion Process (ASEP)

A continuous time Markov process



- Particles move on \mathbb{Z} according to two rules:
 1. A particle waits at x an exponential time with parameter one, and then chooses y with probability $p(x,y)$.
 2. If y is vacant at that time it moves to y , while if y is occupied it remains at x .
- “*Simple*” refers to the fact that jumps are allowed only one step to either the right or left
- “*Asymmetric*” refers to the case $p \neq q$.

Transition Probability: $P_Y(x;t)$

For one particle the probability that the particle is initially at y is at x at time t is

$$P_y(x; t) = \frac{1}{2\pi i} \int_{\mathcal{C}_r} \xi^{x-y-1} e^{t\varepsilon(\xi)} d\xi$$

where

$$\varepsilon(\xi) = \frac{p}{\xi} + q\xi - 1$$

and \mathcal{C}_r is a circle of radius r centered at the origin.

This result is elementary but the generalization to more than one particle is rather subtle

N-particle ASEP

Initial configuration: $Y := \{y_1, y_2, \dots, y_N\}$ with $y_1 < y_2 < \dots < y_N$.

Final configuration: $X := \{x_1, x_2, \dots, x_N\}$ with $x_1 < x_2 < \dots < x_N$.

Let \mathfrak{S}_N denote the permutation group and set

$$U(\xi, \xi') = \frac{p + q\xi\xi' - \xi}{\xi' - \xi}$$
$$A_\sigma(\xi) = \prod_{1 \leq i < j \leq N} \frac{U(\xi_{\sigma(i)}, \xi_{\sigma(j)})}{U(\xi_i, \xi_j)}, \quad \sigma \in \mathfrak{S}_N.$$

Theorem (TW, 2008).

$$P_Y(X; t) = \sum_{\sigma \in \mathfrak{S}_N} \int_{\mathcal{C}_r} \cdots \int_{\mathcal{C}_r} A_\sigma(\xi) \prod_{j=1}^N \xi_{\sigma(j)}^{x_j - y_{\sigma(j)} - 1} e^{t\varepsilon(\xi_j)} d\xi_1 \cdots d\xi_N$$

where \mathcal{C}_r has radius so small that all the poles of A_σ lie outside of \mathcal{C}_r .

Remarks:

- $P_Y(X; t)$ satisfies $P_Y(x; 0) = \delta_{X, Y}$.
- This is a sum of $N!$ terms with each term an N -dimensional contour integral.
- We are ultimately interested in $N \rightarrow \infty$. Not at all clear how to proceed!

- To extract information from $P_Y(x;t)$, we start by looking at *marginal distributions*; the simplest are one-point functions:

$$\mathbb{P}_Y (x_m(t) = x)$$

Must sum $P_Y(X;t)$ over all configurations satisfying $x_m(t) = x$.

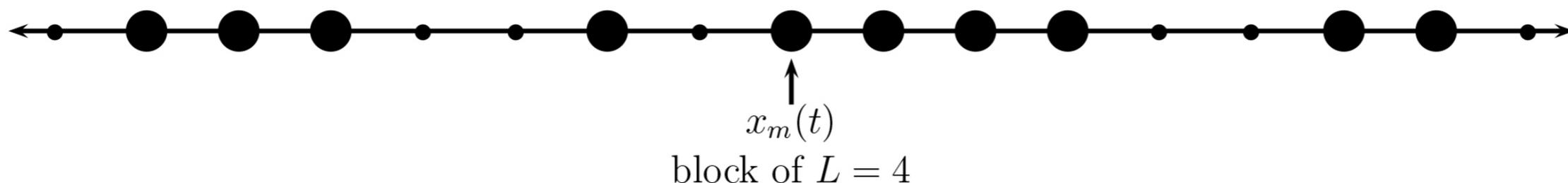
For example, for $m = 2$ we must sum over configurations X

$$X = \{x - v_1, x, x + v_2, x + v_2 + v_3, \dots, x + v_2 + v_3 + \dots + v_N\}$$

where $v_i = 1, 2, 3, \dots$

Second Example: ASEP Blocks

m^{th} particle is the left-most one in a contiguous block of L particles



Case $m=1$, left-most particle

Identity One

Identity One. For $N \geq L$,

$$\begin{aligned} \sum_{\sigma \in \mathfrak{S}_N} \prod_{1 \leq i, j \leq N} U(\xi_{\sigma(i)}, \xi_{\sigma(j)}) \frac{\xi_{\sigma(2)} \xi_{\sigma(3)}^2 \cdots \xi_{\sigma(N)}^{N-1}}{(1 - \xi_{\sigma(L+1)}) \cdots \xi_{\sigma(N)} \cdots (1 - \xi_{\sigma(N-1)} \xi_{\sigma(N)}) (1 - \xi_{\sigma(N)})} \\ = p^{N(N-1)/2} \frac{\mathfrak{f}_L(\xi)}{\prod_i (1 - \xi_i)} \end{aligned}$$

where $\mathfrak{f}_L(\xi)$ are *symmetric polynomials* in the variables $\xi = (\xi_1, \dots, \xi_N)$.

For the definition of $\mathfrak{f}_L(\xi)$ we first define

$$\varphi_L(z_1, \dots, z_L; \xi) = \frac{\prod_{1 \leq j \leq N} U(z_1, \xi_j) U(z_2, \xi_j) \cdots U(z_L, \xi_j)}{z_1^L (qz_1 - p) z_2^{L-1} (qz_2 - p) \cdots z_L (qz_L - p)} \prod_{1 \leq i < j \leq L} \frac{1}{U(z_j, z_i)}$$

$$\text{then } \mathfrak{f}_L(\xi) = p^{L(L+1)/2 - LN} \prod_i \xi_i^L \int_{\Gamma_\xi} \cdots \int_{\Gamma_\xi} \varphi_L(z_1, \dots, z_L; \xi) dz_1 \cdots dz_L,$$

Γ_ξ consists of simple closed curves enclosing the points ξ_j but no other singularities of the integrand.

For $L = 1$,

$$\mathfrak{f}_1(\xi) = 1 - \prod_i \xi_i.$$

but the complexity of \mathfrak{f}_L increases with L .

General m Identity Two

Notation:

- S is a subset of $\{1, 2, \dots, N\}$.
- $\widehat{\xi}_S$ denotes the variables ξ_k with $k \notin S$.
- Set $\tau := p/q < 1$ and recall the τ -binomial coefficients

$$\begin{bmatrix} n \\ k \end{bmatrix}_\tau = \frac{(1 - \tau^n) \cdots (1 - \tau^{n-k+1})}{(1 - \tau) \cdots (1 - \tau^k)}$$

Identity Two: For $0 \leq m \leq N - L$,

$$\sum_{|S|=m} \prod_{\substack{i \in S \\ j \notin S}} U(\xi_i, \xi_j) \cdot \mathfrak{f}_L(\widehat{\xi}_S) = q^{m(N-m)} \begin{bmatrix} N - L \\ m \end{bmatrix}_\tau \mathfrak{f}_L(\xi)$$

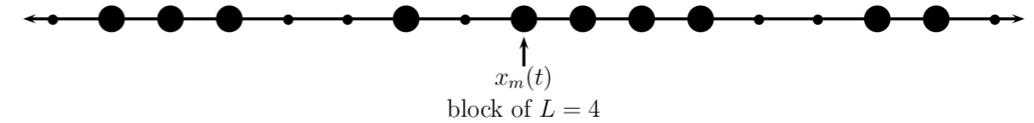
where $\mathfrak{f}_L(\xi)$ are the symmetric polynomials from *Identity One* and

$$U(\xi, \xi') = \frac{p + q\xi\xi' - \xi}{\xi' - \xi}.$$

What do the Identities buy for you?

Notation:

- $\mathcal{P}_{L,Y}(x, m, t)$: probability that at time t the m th particle from the left is the beginning of a block of length L starting at x .



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$$I_L(x, Y, \xi) := \prod_{1 \leq i < j \leq N} \frac{1}{U(\xi_i, \xi_j)} \prod_i \frac{1}{1 - \xi_i} f_L(\xi) \prod_i \left(\xi_i^{x-y_i-1} e^{\varepsilon(\xi_i)t} \right)$$

- S a subset of $\{1, \dots, N\}$, S^c complement of S .
- $I_L(x, Y_S, \xi_S)$ indices lie in S .
- $\sigma(S^c)$ is the sum of the elements in S^c .

Theorem (TW, $L = 1$, 2008; general L , 2017): For $q > 0$

$$\begin{aligned} \mathcal{P}_{L,Y}(x, m, t) &= p^{(N-m+1)(N-m)/2} q^{(m-1)(N-m/2)} \sum_{|S^c| < m} (-1)^{m-1-|S^c|} \begin{bmatrix} |S| - L \\ m - 1 - |S^c| \end{bmatrix}_\tau \\ &\quad \times \frac{q^{\sigma(S^c) - N|S^c|}}{p^{\sigma(S^c) - |S^c|(|S^c|+1)/2}} \int_{\mathcal{C}_r} \cdots \int_{\mathcal{C}_r} I_L(x, Y_S, \xi_S) d^{|S|} \xi \end{aligned}$$

Remarks:

- The proof for general L proceeds exactly the same as for $L = 1$ given the general L identities and the fact that $f_L(\xi)$ are polynomials—no new poles introduced in the argument.
- As was the case for $L = 1$, there is a formula for $\mathcal{P}_{L,Y}(x, m, t)$ but with integrations over large contours. In this expression one can let $N \rightarrow \infty$.

Large contour representation

Notation:

- $\mathcal{P}_{L,Y}(x, m, t)$: probability that at time t the m th particle from the left is the beginning of a block of length L starting at x .

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$$I_L(x, Y, \xi) := \prod_{1 \leq i < j \leq N} \frac{1}{U(\xi_i, \xi_j)} \prod_i \frac{1}{1 - \xi_i} f_L(\xi) \prod_i \left(\xi_i^{x-y_i-1} e^{\varepsilon(\xi_i)t} \right)$$

- S a subset of $\{1, \dots, N\}$.
- $I_L(x, Y_S, \xi_S)$ indices lie in S .
- $\sigma(S)$ is the sum of the elements in S .

Theorem (TW, $L = 1$, 2008; general L , 2017): For $q > 0$

$$\begin{aligned} \mathcal{P}_{L,Y}(x, m, t) &= (-1)^{m+1} p^{m(m-1)/2} \sum_{|S| \geq m+L-1} q^{(m-1)(|S|-m/2)} \begin{bmatrix} |S| - L \\ m - 1 \end{bmatrix}_\tau \\ &\quad \times \frac{p^{\sigma(S)-m|S|}}{q^{\sigma(S)-|S|(|S|+1)/2}} \int_{\mathcal{C}_R} \cdots \int_{\mathcal{C}_R} I_L(x, Y_S, \xi_S) d^{|S|} \xi \end{aligned}$$

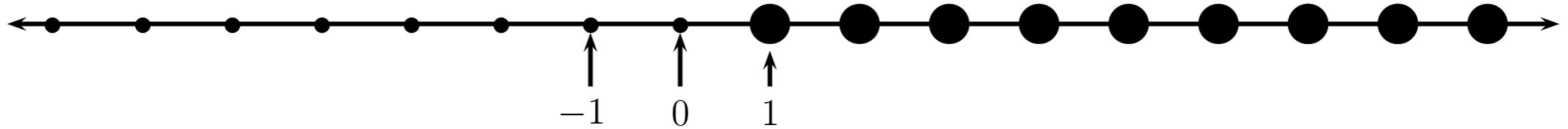
where R is so large that the poles of the integrand lie inside \mathcal{C}_R .

Remarks:

- This theorem extends to **infinite systems unbounded on the right**. The sum is then taken over finite subsets of \mathbb{Z}^+ .
- Up to this point the initial configuration $Y = \{y_1, y_2, \dots\}$, $y_1 < y_2 < \dots$, is completely general (though bounded below). We now turn to the special case of *step initial condition*.

Step Initial Condition

Drift to the left, $p < q$
Particles initially occupy \mathbb{Z}^+



Remarks:

- In the stochastic growth interpretation of ASEP, the step initial condition corresponds to the *droplet initial condition*.
- We are interested in $\mathcal{P}_{L, \mathbb{Z}^+}(x, m, t)$.
- One starts with the *large contour representation* of $\mathcal{P}_{L, \mathbb{Z}^+}(x, m, t)$, and then first sums over all S with $|S|$ equal to a fixed k .

Fredholm Determinant Representation

Notation:

- Denote by $K_{L,x}(z)$ the integral operator acting on functions on \mathcal{C}_R with kernel

$$K_{L,x}(\xi, \xi'; z) = K_x(\xi, \xi') \prod_{j=1}^L U(z_j, \xi), \text{ where}$$

$$K_x(\xi, \xi') = \frac{\xi^x e^{\varepsilon(\xi)t}}{p + q\xi\xi' - \xi}$$

- τ -Pochhammer symbol, $(\lambda; \tau)_m := \prod_{j=0}^{m-1} (1 - \lambda\tau^j)$.

Theorem (TW, $L = 1$, 2008; general L , 2017). For $p, q > 0$,

$$\mathcal{P}_{L, \mathbb{Z}^+}(x, m, t) = (-1)^{L-1} p^{L(L+1)/2} \tau^{-(m-1)(L-1)}$$

$$\times \int_{\Gamma_{0,\tau}} \cdots \int_{\Gamma_{0,\tau}} \frac{1}{z_1^L (qz_1 - p) z_2^{L-1} (qz_2 - p) \cdots z_L (qz_L - p)} \prod_{i < j} \frac{1}{U(z_j, z_i)}$$

$$\times \left[\int \frac{\det(I - p^{-L} q \lambda K_{L, x+L-1}(z))}{(\lambda; \tau)_m} \frac{d\lambda}{\lambda^L} \right] dz_L \cdots dz_1.$$

Remarks:

- The z -iterated integral is interpreted as follows: First take the sum of the residues at $z_L = 0$ and $z_L = \tau$. In the resulting integrand take the sum of the residues at $z_{L-1} = 0$ and $z_{L-1} = \tau$; and so on.
- The λ -integration is over a contour enclosing the singularities of the integrand at τ^{-j} for $j = 0, \dots, m-1$.
- For $L = 1$, evaluating the z_1 -integral leads to the result

$$\mathbb{P}_{\mathbb{Z}^+}(x_m(t) \leq x) = \int \frac{\det(I - q\lambda K_x)}{(\lambda; \tau)_m} \frac{d\lambda}{\lambda}$$

which is the 2008 result.

*But what does all this ASEP stuff
have to do with
Painlevé?*

A Limit Theorem

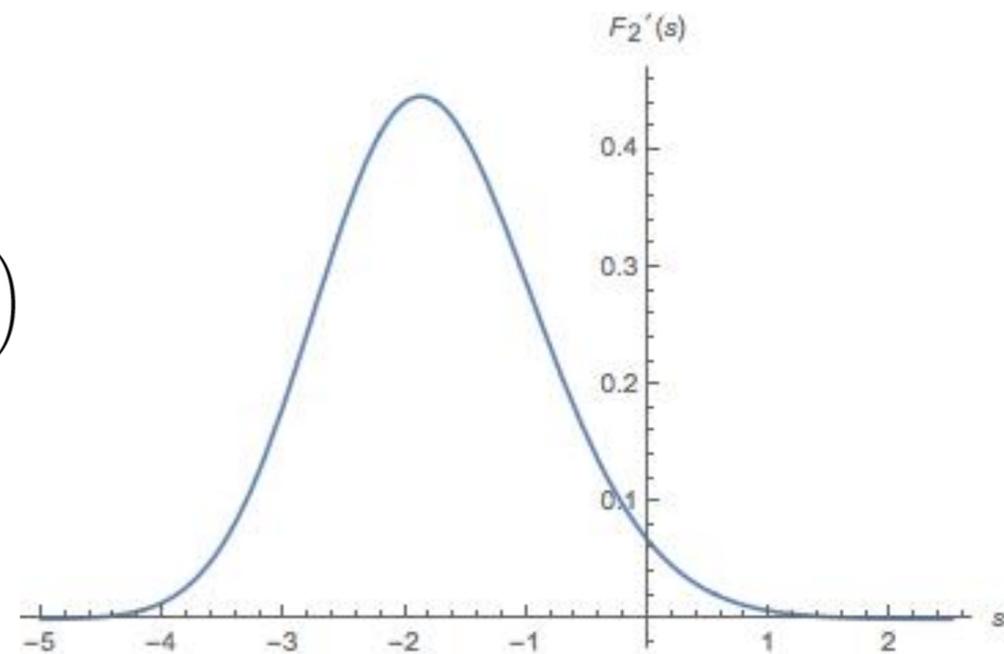
Notation:

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$$F_2(s) := \exp \left(- \int_s^\infty (x - s) q^2(x) dx \right)$$

where q is the Hastings-McLeod solution of P_{II}

$$\frac{d^2 q}{dx^2} = xq + 2q^3.$$



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$$\sigma = \frac{m}{t}, \quad c_1 = -1 + 2\sqrt{\sigma}, \quad c_2 = \sigma^{-1/6}(1 - \sqrt{\sigma})^{2/3}, \quad \gamma = q - p,$$

Theorem (TW, $L = 1$, 2009). When $0 \leq p < q$ we have

$$\lim_{t \rightarrow \infty} \mathbb{P} \left(\frac{x_m(t/\gamma) - c_1 t}{c_2 t^{1/3}} \leq s \right) = F_2(s)$$

uniformly for σ in a compact subset of $(0, 1)$.

Remarks:

- For $p = 0$ (TASEP), this result was first proved by Johansson (2000).
- The above Theorem was the first “KPZ Universality Theorem” outside the class of *determinantal processes*. (TASEP is a determinantal process.)
- The proof required a deformation of the K_x -operator to a new J -operator that (essentially) left the Fredholm determinant of K_x unchanged. The Fredholm determinant of the J -operator is amenable to asymptotic analysis.
- The J -operator representation was the starting point for the Amir-Corwin-Quastel and (independently) the Sasamoto-Spohn analysis of the KPZ crossover scaling function.

Future Work

- Extend the *J*-kernel analysis to general L . (Essentially done.)
- Carry out an asymptotic analysis in the KPZ limit. (Work in progress.)
- Develop a better understanding of the identities. For $L=1$ a deeper interpretation was found by Borodin, Corwin, Petrov and Sasamoto (Commun. Math. Phys. 329 (2015), 1167—1245).
- “ASEP Gaps” (work in progress)

*Thank you
for your attention*